

Chapter 7 problems

LAST NAME

FIRST NAME

Problem 7.1

(a) Write a computer program to calculate the chemical potential for n non-interacting electrons per unit volume at temperature, T .

(b) Calculate the value of the chemical potential for the case when electrons of effective mass $m^* = 0.07 \times m_0$ and carrier density $n = 1.5 \times 10^{18} \text{ cm}^{-3}$ are at temperature $T = 300 \text{ K}$

(c) Repeat (b), only now for the case when electrons have effective electron mass $m^* = 0.50 \times m_0$.

(d) Plot the Fermi-Dirac distribution function for the situations described by (b) and (c).

(e) Repeat (b), (c), and (d), only now for the case when temperature $T = 77 \text{ K}$.

Your answer should include a print out of your computer program and plots.

Problem 7.2

(a) Show that
$$\frac{1}{e^{(E-\mu)/k_B T} + 1} = 1 - \frac{1}{e^{(\mu-E)/k_B T} + 1}$$

(b) A semiconductor consists of a valance band with electron energy dispersion relation $E_{VB} = E(\mathbf{k})$ and a conduction band with electron energy dispersion relation such that $E_{CB} = E_0 - E(\mathbf{k})$, where E_0 is a constant such that the conduction band and valance band are separated by an energy band gap, E_g . Show that when particle number is conserved, the chemical potential is in the middle of the band gap with value $\mu = E_0/2$ and is independent of temperature.

Problem 7.3

(a) Calculate the average energy of electrons in a three-dimensional gas of electrons.

Show that in the low temperature limit $\langle E_{3D}(T \rightarrow 0 \text{ K}) \rangle = \frac{3}{5} E_F$ and in the high temperature limit

$\langle E_{3D}(T \rightarrow \infty \text{ K}) \rangle = \frac{3}{2} k_B T$.

(b) Calculate the average energy of electrons in a two-dimensional gas of electrons.

Show that in the low temperature limit $\langle E_{2D}(T \rightarrow 0 \text{ K}) \rangle = \frac{1}{2} E_F$ and in the high temperature limit

$\langle E_{2D}(T \rightarrow \infty \text{ K}) \rangle = k_B T$.

Problem 7.4

The anti-symmetric wave function that describes two identical indistinguishable non-interacting particles is given by the Slater determinant

$$\Psi_a(x_1, x_2) = \frac{1}{\sqrt{2}} \begin{vmatrix} \psi_1(x_1) & \psi_1(x_2) \\ \psi_2(x_1) & \psi_2(x_2) \end{vmatrix}$$

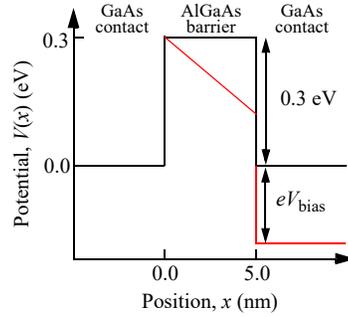
where rows label the single-particle state and columns label the particle. Position coordinate for particle 1 is x_1 and for particle 2 it is x_2 .

(a) Plot $\Psi_a(x_1, x_2)$ for the case when ψ_1 is the single-particle ground state of a one-dimensional rectangular potential well with infinite barrier energy and ψ_2 is the first excited state. Comment on the value of $|\Psi_a(x_1, x_2)|^2$ when $x_1 = x_2$.

(b) Repeat the calculation in (a) but now for the case when ψ_1 is the single-particle first excited state and ψ_2 is the second excited state. Comment on your results.

(c) Repeat the calculations in (a) and (b) for symmetric wave functions $\Psi_s(x_1, x_2)$. Comment on your results.

Problem 7.5



(a) Consider the conduction band minimum potential profile shown in the Fig. It consists of two GaAs contact layers and an AlGaAs potential energy barrier region. The contacts have the same n -type impurity concentration and the AlGaAs is intrinsic. The current due to a single electron in state $|\mathbf{k}\rangle$ with energy E_k is

$$J_e = e \frac{\hbar k_{\perp}}{m} T(E_{\perp}),$$

where $k_{\perp} = (\sqrt{2m(E_{\perp} - V(x))})/\hbar$ is the component of \mathbf{k} perpendicular to the layer interface and $T(E_{\perp})$ is the transmission coefficient. The total current flowing left-to-right involves all electron states in the left contact and so requires integrations over both k_{\perp} and k_{\parallel} , where $k_{\parallel} = \sqrt{2mE_{\parallel}}/\hbar$ is the component of \mathbf{k} parallel to the contact-barrier interface, and $E_k = E_{\perp} + E_{\parallel}$. If the probability of an electron mass m and charge e occupying state $|\mathbf{k}\rangle$ is given by the Fermi function $f(E_k, \mu)$, the current due to the left contact is

$$J_L = e \int T(E_{\perp}) \frac{\hbar k_{\perp} dk_{\perp}}{m 2\pi} \int f(E_k, \mu) \frac{d^2 k_{\parallel}}{(2\pi)^2}$$

where μ is the chemical potential in the contact. Total current is the difference between the left contact and the right contact current. A positive bias voltage, V_{bias} , lowers the chemical potential energy of the right hand contact by eV_{bias} , and the total current is

$$J = e \int T(E_{\perp}) \frac{\hbar k_{\perp} dk_{\perp}}{m 2\pi} \int (f(E_k, \mu) - f(E_k, \mu - eV_{\text{bias}})) \frac{d^2 k_{\parallel}}{(2\pi)^2}$$

Starting from this expression, use the one and two-dimensional densities of states to convert the integrals to energy and evaluate the integration over k_{\parallel} to show that

$$J = \frac{emk_B T}{\pi^2 \hbar^3} \int_0^\infty T(E_\perp) \ln \left(\frac{1 + e^{(\mu - E_\perp)/k_B T}}{1 + e^{(\mu - E_\perp - eV_{\text{bias}})/k_B T}} \right) dE_\perp$$

(b) Use MATLAB to calculate current density through the potential barrier. Use an effective electron mass $m = 0.07 \times m_0$ and an initial voltage bias range of $0 < V_{\text{bias}} < 0.3$ V. The code developed for Problem 4.3 may be used to calculate the transmission probability, $T(E_\perp)$. As in the Fig., assume that the potential change due to V_{bias} appears linearly with position across the barrier. Plot the calculated current density using both linear and log scales. Explain the dependence of current density on voltage bias that you observe when impurity concentration has value $n = 10^{18} \text{ cm}^{-3}$ and $n = 10^{16} \text{ cm}^{-3}$ and when temperature has value $T = 300$ K and $T = 4.2$ K. Explain what you observe when you extend the voltage bias range to $0 < V_{\text{bias}} < 2.3$ V.

Problem 7.6

The minimum value of chemical potential μ_{min} at finite absolute temperature T may be found by assuming a particle distribution function that is obtained in the $T \rightarrow \infty$ limit.

(a) Show that in this case a three-dimensional electron gas of fixed density n_{3D} has minimum chemical potential

$$\mu_{\text{min}}^{3D}(T) = k_B T \ln \left(\frac{n_{3D}}{2} \left(\frac{2\pi\hbar^2}{m_0 k_B T} \right)^{\frac{3}{2}} \right)$$

where m_0 is the electron mass.

(b) Find the expression for $\mu_{\text{min}}^{2D}(T)$ of a two-dimensional electron gas.

(c) Find the expression for $\mu_{\text{min}}^{1D}(T)$ of a one-dimensional electron gas.

(d) Plot the exact μ and minimum μ_{min} from part (a), (b), and (c), for normalized chemical potential μ/E_F as a function of normalized thermal energy $k_B T/E_F$, where E_F is the Fermi energy. Explain the differences you observe.

Problem 7.7

Three particles in a one-dimensional harmonic oscillator potential obey the Pauli exclusion principle. Assuming that any two microscopically distinguishable arrangements of the system with the same total energy are equally likely (the ergodic theorem), plot the probability of occupation as a function of energy for (a) the lowest energy state of the system, (b) when the total energy is $E_{\text{total}} = 13.5 \times \hbar\omega$, and (c) when total energy is $E_{\text{total}} = 48.5 \times \hbar\omega$. (d) Repeat (a) - (c) for the case when there are five particles. Comment on anything you learn.

Problem 7.8

The Hubbard model is used to describe systems in terms of interacting particles on a lattice. For a one-dimensional (1D) lattice with only two sites, the Hamiltonian for electrons of spin $\sigma = \uparrow$ or \downarrow is

$$\hat{H} = -t \sum_{\langle i,j \rangle, \sigma} (\hat{c}_{2,\sigma}^\dagger \hat{c}_{1,\sigma} + \hat{c}_{1,\sigma}^\dagger \hat{c}_{2,\sigma}) + U \sum_i \hat{n}_{i,\uparrow} \hat{n}_{i,\downarrow}$$

where t is the nearest neighbor hopping matrix element, U is the additional energy associated with two electrons of opposite spin occupying the same lattice site, i , the electron creation and annihilation operators are \hat{c}^\dagger and \hat{c} respectively, and \hat{n} is the electron number operator.

(a) If there are a total of two electrons in the system ($N = 2$), write down all possible states in the Fock particle number basis.

(b) Using the Fock basis and \hat{H} , calculate all matrix elements and derive expression for the eigenenergies in terms of the coupling strengths U and t by finding the eigenvalues for the matrix. Plot the eigenenergies as a function of the ratio U/t .

(c) Repeat (a) and (b) for the same one-dimensional lattice with two sites, but now for the cases when $N = 1, 3, 4$ and comment on what you learn.

Problem 7.9

Explain why the statement “The wave function $\psi(\mathbf{r})$ is a function of three-dimensional real-space” is incorrect.

Problem 7.10

An electron gas obeys Fermi-Dirac statistics at temperature T and is constrained to motion in the x -direction. A potential barrier of infinite energy is placed at position $x = 0$ such that potential $V(x > 0) = \infty$ and $V(x \leq 0) = 0$.

(a) By considering a region that confines electrons from $x = 0$ to $x = L < 0$, find the normalized electron energy eigenstates and explain how taking the limit $L \rightarrow -\infty$ creates a continuum of states.

(b) At sufficiently high temperatures the Fermi-Dirac distribution may be approximated by a Maxwell-Boltzmann distribution. Find an analytic expression for the electron density $n = n(x, T)$ in terms of a characteristic length scale, $\lambda = \hbar / \sqrt{2m_0 k_B T}$. Explain why n has a smaller value near $x = 0$ than when $x \ll -\lambda$. Plot the density $n = n(x, T)$ at temperatures $T = 10, 100,$ and 300 K for electrons with effective electron mass $m^* = 0.07 \times m_0$.

(c) When temperature $T < E_F / k_B$ the Fermi-Dirac distribution may not be approximated by a Maxwell-Boltzmann distribution. Find and explain the spatial dependence of carrier concentration $n = n(x, T)$ for temperatures $T = 300$ K, $T = 77$ K, and $T = 4$ K.

You may wish to make use for the standard integrals $\int_0^\infty e^{-a\xi^2} d\xi = \frac{1}{2} \sqrt{\frac{\pi}{a}}$ and

$$\int_{\xi=0}^{\xi=\infty} e^{-a\xi^2} \cos(b\xi) d\xi = \sqrt{\frac{\pi}{4a}} e^{-\frac{b^2}{4a}} \text{ for } \text{Re } a > 0.$$

Problem 7.11

At temperatures $T = 3 \text{ K}$ and $T = 300 \text{ K}$ a fermion system has chemical potential $\mu_{T=3 \text{ K}} = 52.1 \text{ meV}$ and $\mu_{T=300 \text{ K}} = 38.9 \text{ meV}$ respectively. Use MATLAB to plot the Fermi-Dirac distribution $f(E)$ and the Bose distribution $g(E)$ as a function of energy $0 < E < 200 \text{ meV}$. Explain the difference in the curves you obtain. What do you expect to happen to the distribution functions in the high-temperature limit?

Problem 7.12

(a) Show that transmission and reflection amplitude of a single photon at a perfect lossless symmetric 50:50 beam splitter is $t = i/\sqrt{2}$ and $r = -1/\sqrt{2}$, respectively.

(b) The total number of identical indistinguishable photons $n_{\text{tot}} = n_1 + n_2$ at the input to a perfect lossless symmetric 50:50 beam splitter consists of integer n_1 photons at input port 1 and integer n_2 photons at input port 2. In general, the quantum amplitude of integer n_3 and integer n_4 indistinguishable photons appearing at output ports 3 and 4 respectively of the beam splitter is

$$|n_1, n_2, n_3, n_4\rangle = (-1)^{n_1} \left(\frac{1}{2}\right)^{\frac{n_1+n_2}{2}} \sum_k (-1)^k \sqrt{\binom{n_1}{k} \binom{n_2}{n_3-k} \binom{n_3}{k} \binom{n_4}{n_1-k}}$$

where, because the total number of particles is conserved, $n_4 = n_1 + n_2 - n_3$. In this expression, the binomial coefficient

$$\binom{n}{k} = \frac{n!}{k!(n-k)!}$$

If k is negative or greater than n the binomial coefficient should be set to zero.

Write a computer program to calculate the probability of all possible input-output combinations when $n_{\text{tot}} = 8$. What is the probability distribution of detecting n_3 photons ($n_3 = 0, 1, 2, \dots, 8$) when $n_1 = 4$ and $n_2 = 8$?

Calculate the probability of all possible input-output combinations when $n_{\text{tot}} = 128$. Comment on what you learn.

Problem 7.13

The anti-symmetric state ψ_a describing N identical indistinguishable non-interacting fermion particles may be found using the Slater determinant (Eq. 7.18)

$$\Psi_a(x_1, x_2, \dots, x_N) = \frac{1}{\sqrt{N!}} \begin{vmatrix} \Psi_1(x_1) & \Psi_1(x_2) & \dots & \Psi_1(x_N) \\ \Psi_2(x_1) & \Psi_2(x_2) & \dots & \Psi_2(x_N) \\ \cdot & \cdot & \dots & \cdot \\ \cdot & \cdot & \dots & \cdot \\ \Psi_N(x_1) & \dots & \dots & \Psi_N(x_N) \end{vmatrix}$$

in which the rows label the i -th particle state and columns label j -th particle at position x_j . The anti-symmetric state is

$$\Psi_a(x_1, x_2, x_3, \dots, x_N) = \langle x_1, x_2, x_3, \dots, x_N | n_1, n_2, n_3, \dots, n_N \rangle$$

where the n_k in the particle-number basis $|n_1, n_2, n_3, \dots, n_N\rangle$ can only have a value of 1 or 0 as required by the Pauli exclusion principle. The interchange of any two particles causes the sign of the multi-particle wave function ψ_a to change, since it involves the interchange of two columns in the Slater determinant.

(a) Use the Slater determinant to show that the fermion annihilation \hat{c}_k acting on the many-electron state gives

$$\hat{c}_k |n_1, \dots, n_k = 1, \dots, n_N\rangle = (-1)^{\sum_{j < k} n_j} |n_1, \dots, n_k = 0, \dots, n_N\rangle$$

when $n_k = 1$ and

$$\hat{c}_k |n_1, \dots, n_k = 0, \dots, n_N\rangle = 0$$

when $n_k = 0$.

Show that the creation operator \hat{c}_k^\dagger acting on the many-electron state gives

$$\hat{c}_k^\dagger |n_1, \dots, n_k = 1, \dots, n_N\rangle = 0$$

when $n_k = 1$ and

$$\hat{c}_k^\dagger |n_1, \dots, n_k = 0, \dots, n_N\rangle = (-1)^{\sum_{j < k} n_j} |n_1, \dots, n_k = 1, \dots, n_N\rangle$$

when $n_k = 0$.

(b) Use part (a) to obtain the anti-commutation relations

$$\{\hat{c}_j, \hat{c}_i^\dagger\} = \hat{c}_j \hat{c}_i^\dagger + \hat{c}_i^\dagger \hat{c}_j = \delta_{ji}$$

$$\{\hat{c}_j, \hat{c}_i\} = 0$$

$$\{\hat{c}_j^\dagger, \hat{c}_i^\dagger\} = 0$$
